

Obstacles for Magnetic Hamiltonians: SSF near Landau levels

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with G. Raikov

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In honor of André Martinez

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Schrödinger operators with constant magnetic field $B = (0, 0, b)$:

$$H_0 := -(\nabla^A)^2 = \left(D_1 + \frac{b}{2}x_2\right)^2 + \left(D_2 - \frac{b}{2}x_1\right)^2 + D_3^2$$

$$D_j := -i\frac{\partial}{\partial x_j}, \quad \nabla_j^A := \nabla_{x_j} - iA_j, \quad A = \left(\frac{b}{2}x_2, -\frac{b}{2}x_1, 0\right)$$

H_0 is s.a. on the Magnetic Sobolev spaces $\mathfrak{D}(H_0) := H_A^2(\mathbb{R}^3)$

$$\|u\|_{H_A^s(\Omega)}^2 := \sum_{\alpha \in \mathbb{Z}_+^3: 0 \leq |\alpha| \leq s} \int_{\Omega} |(\nabla^A)^\alpha u|^2 dx.$$

$\Omega_{\text{in}} \subset \mathbb{R}^3$: a connected regular bounded domain. $\Omega_{\text{ex}} := \mathbb{R}^3 \setminus \Omega_{\text{in}}$,

$\Gamma := \partial\Omega_{\text{in}} = \partial\Omega_{\text{ex}}$, ν : outward normal vector at Γ (w.r.t. Ω_{in})

Dirichlet realization on Ω_j : H_{+j} , $j = \text{ex, in}$, restriction of H_0 on

$$\mathfrak{D}(H_{+j}) := \{u \in H_A^2(\Omega_j) \mid u|_{\Gamma} = 0\}.$$

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Free Magnetic Schrödinger operators in \mathbb{R}^3 :

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We consider geometric perturb.: in $L^2(\mathbb{R}^3) = L^2(\Omega_{\text{in}}) \oplus L^2(\Omega_{\text{ex}})$,

$$H_{\pm} := H_{\pm,\text{in}} \oplus H_{\pm,\text{ex}}.$$

Boundary condition on Γ : " + ": Dirichlet " - ": Neumann

We have:

$$H_- \leq H_0 \leq H_+$$

$$\sigma(H_{\pm,\text{in}}) = \sigma_{\text{disc}}(H_{\pm,\text{in}}) \subset (0, +\infty)$$

$$\sigma_{\text{ess}}(H_{\pm,\text{ex}}) = \sigma_{\text{ess}}(H_0) = \sigma(H_0) = [b, +\infty)$$

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Recall: spectrum of H_0 :

- The Landau Hamiltonian, on $L^2(\mathbb{R}^2)$:

$$H_{0,\perp} = \left(i \frac{\partial}{\partial x_1} - \frac{bx_2}{2} \right)^2 + \left(i \frac{\partial}{\partial x_2} + \frac{bx_1}{2} \right)^2 \approx (D_{y_1}^2 + y_1^2) \otimes I_{y_2}$$

Spectrum: $\begin{array}{cccccc} & b & & 3b & & 5b & & 7b & & 9b \\ & | & & | & & | & & | & & | \end{array}$

$$\sigma(H_{0,\perp}) = \sigma_{\text{ess}}(H_{0,\perp}) = \{\Lambda_j = b(2j+1); j \in \mathbb{N}\},$$

$\Lambda_j = b(2j+1)$: **Landau Level**, $\text{mult.}(\Lambda_j) = \infty$

- The 3D Magnetic Schrödinger operator, on $L^2(\mathbb{R}^3)$:

$$H_0 = H_{0,\perp} - \frac{\partial^2}{\partial x_3^2}$$

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Goal: measure of the influence of the obstacle on the spectrum
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The Spectral Shift Function (SSF) for the pair (H_{\pm}, H_0) :
In the sense of the distributions

$$\xi(E; H_{\pm}, H_0) := -\mathrm{Tr} \left(\mathbf{1}_{(-\infty, E)}(H_{\pm}) - \mathbf{1}_{(-\infty, E)}(H_0) \right)$$

For a.e. $E \in [b, \infty) = \sigma_{\mathrm{ac}}(H_0)$, the Birman-Krein formula implies

$$\det S(E; H_{\pm}, H_0) = e^{-2\pi i \xi(E; H_{\pm}, H_0)}$$

where $S(E; H_{\pm}, H_0)$ is the scattering matrix

Proposition

The SSF coincide a.e. with $\tilde{\xi}(\cdot; H_{\pm}, H_0)$ which is :

- bounded on every compact subset of $(0, \infty) \setminus b(2\mathbb{Z}_+ + 1)$
- continuous on $(0, \infty) \setminus (\sigma_p(H_{\pm}) \cup b(2\mathbb{Z}_+ + 1))$

Question 1 Behavior of $\xi(E; H_{\pm}, H_0)$ as $E \rightarrow \Lambda_q = b(2q + 1)$?

Question 2 How does this behavior depend on the obstacle?

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II) Singularities of the SSF near Λ_q :

Remark: $\xi(E; H_{\pm}, H_0) := -\text{Tr} \left(\mathbf{1}_{(-\infty, E)}(H_{\pm}) - \mathbf{1}_{(-\infty, E)}(H_0) \right)$

$$= -\text{Tr} \left(\mathbf{1}_{(-\infty, E)}(H_{\pm, \text{in}}) \oplus \mathbf{1}_{(-\infty, E)}(H_{\pm, \text{ex}}) - \mathbf{1}_{(-\infty, E)}(H_0) \right)$$
$$= -\text{Tr} \left(\mathbf{0} \oplus \mathbf{1}_{(-\infty, E)}(H_{\pm, \text{ex}}) - \mathbf{1}_{(-\infty, E)}(H_0) \right) - \text{Tr} \left(\mathbf{1}_{(-\infty, E)}(H_{\pm, \text{in}}) \right)$$

Since the spectrum of $H_{\pm, \text{in}}$ is discret, near each Λ_q

$$\text{Tr} \left(\mathbf{1}_{(-\infty, E)}(H_{\pm, \text{in}}) \right) = O(1)$$

Then as $E \rightarrow \Lambda_q$

$$\xi(E; H_{\pm}, H_0) = \xi(E; H_{\pm, \text{ex}}, H_0) + O(1)$$

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For $\lambda \in (0, e^{-1})$ set $\ln_2(\lambda) := \ln |\ln \lambda|$, $\Phi_0(\lambda) := \frac{|\ln \lambda|}{\ln_2(\lambda)}$

Theorem 1

Let $q \in \mathbb{Z}_+$. Then we have

$$\xi(\Lambda_q - \lambda; H_+, H_0) = O(1),$$

$$\xi(\Lambda_q - \lambda; H_-, H_0) = -\frac{1}{2}\Phi_0(\lambda)(1 + o(1)),$$

$$\xi(\Lambda_q + \lambda; H_{\pm}, H_0) = \pm\frac{1}{4}\Phi_0(\lambda)(1 + o(1)),$$

as $\lambda \downarrow 0$.

Remark: The main contribution, $\Phi_0(\lambda)$, is independent of Ω_{in} .

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$$H_{0,\perp} = (D_1 - \frac{bx_2}{2})^2 + (D_2 + \frac{bx_1}{2})^2, \quad H_{\perp} = H_{0,\perp} + V \text{ or } H_{\perp,\pm}$$

If $V > 0$ or $H_{\perp,+}$: infinitely many eig. **above** each Landau level

If $V < 0$ or $H_{\perp,-}$: infinitely many eig. **below** each Landau level



On $\mathbb{R}^2 \setminus \mathcal{O}_{\text{in}}$ or for $V = \pm \mathbf{1}_{\mathcal{O}_{\text{in}}}$, $\mathcal{O}_{\text{in}} \subset \mathbb{R}^2$ bounded. Near Λ_q :

$$\mathcal{N}_{\pm}(\lambda) := \#\{\text{eig.}(H_{\perp}) \in \Lambda_q \pm (\lambda, r_0)\}, \quad 0 < \lambda < r_0 < 2b,$$

Theorem (RaWa02, PuRo07, Pe07, GoKaPe14)

$$\mathcal{N}_{\pm}(\lambda) \sim \Phi_0(\lambda) := |\ln \lambda| (\ln |\ln \lambda|)^{-1}, \quad \text{as } \lambda \searrow 0$$

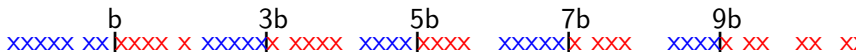
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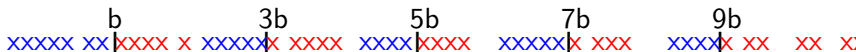
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Known results : 3D case

H : Perturbation of H_0 by

- potentials ($V \rightarrow 0$ at ∞):
Av-He-Si '78, Fe-Ra'04, Bo-B.-Ra '07, '10
- by obstacles: B-Sa'16

The perturbation may generate :

- Accumulation of discrete eigenvalues to the first Landau Level
- Singularities of the Spectral Shift Function (Scattering Phase) at the Landau Levels: $SSF(\Lambda_q \pm \lambda) \rightarrow \infty$ as $\lambda \searrow 0$.
- Accumulation of resonances or embedded eigenvalues to Λ_q

$$H_V = H_0 + V, \quad V \longrightarrow 0 \text{ at } \infty$$

1) EIGENVALUES:

For $V \leq 0$, $V(X_\perp, x_3) = V(|X_\perp|, x_3)$, $X_\perp = (x_1, x_2)$
[Avron-Herbst-Simon '78]

$$-2b < V(X_\perp, x_3) \leq -C \mathbf{1}_K(X_\perp) v(x_3), \quad K \subset \mathbb{R}^2, \quad C > 0$$

- * If $v(x_3) = \mathbf{1}_{\tilde{K}}(x_3)$: at least one eig. in $(\Lambda_{q-1}, \Lambda_q)$, $q \in \mathbb{N}$
- * If $v(x_3) = \langle x_3 \rangle^{-\alpha}$, $\alpha \in (0, 2)$: $\exists (\lambda_{q,j})_{q,j}$ eig. of H_V , $\lambda_{q,j} \rightarrow \Lambda_q$



For $0 \leq V \leq C \langle X_\perp \rangle^{-m_\perp} \langle x_3 \rangle^{-m_3}$, $m_\perp > 0$, $m_3 > 2$ only a
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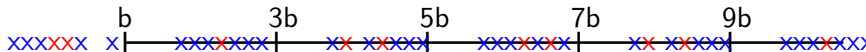
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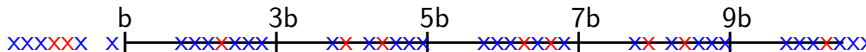
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Compact. supp. pertub. of H_0 ($H_{\pm} = H_0 \pm V$, $V \geq 0$ or H_{\pm} with B.C.)

2) SSF: (Fe-Ra'04, B-Ra'19)

Assume $\frac{1}{C} \mathbf{1}_{K_0} \leq V \leq C \mathbf{1}_{K_1}$, $K_0 \subset K_1 \subset \mathbb{R}^3$ bounded

Then, near a fixed Landau level Λ_q :

- For the pair (H_-, H_0) , as $\lambda \downarrow 0$,

$$\xi(\Lambda_q - \lambda) \sim -\frac{|\ln \lambda|}{2 \ln |\ln \lambda|} \quad \Lambda_q \quad \xi(\Lambda_q + \lambda) \sim \frac{|\ln \lambda|}{4 \ln |\ln \lambda|}$$

- For the pair (H_+, H_0) , as $\lambda \downarrow 0$,

$$\xi(\Lambda_q - \lambda) = O(1) \quad \Lambda_q \quad \xi(\Lambda_q + \lambda) \sim \frac{|\ln \lambda|}{4 \ln |\ln \lambda|}$$

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3) RESONANCES (Bo-B.-Ra '07, '10, B-Sa '16)

$$\#\{z \text{ Res.}; \text{dist}(z, \Lambda_q) \geq r\} \sim \frac{|\ln r|}{\ln |\ln r|}, \quad r \downarrow 0,$$

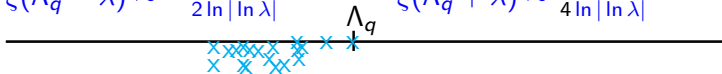
Compact. supp. pertub. of H_0 ($H_{\pm} = H_0 \pm V$, $V \geq 0$ or H_{\pm} with B.C.)

Then, near a fixed Landau level Λ_q :

- For the pair (H_-, H_0) , as $\lambda \downarrow 0$,

$$\xi(\Lambda_q - \lambda) \sim -\frac{|\ln \lambda|}{2 \ln |\ln \lambda|}$$

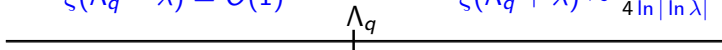
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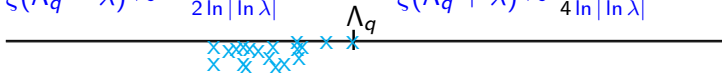
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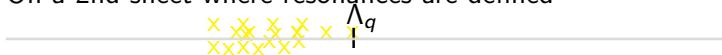
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On a 2nd sheet where resonances are defined



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$$\#\{z \text{ Res.}; \text{dist}(z, \Lambda_q) \geq r\} \sim \frac{|\ln r|}{\ln |\ln r|}, \quad r \downarrow 0,$$

Dependence on the domain in the 2D case

Comp. Supp. perturbations of the Landau Ham. $H_{0,\perp}$

Filonov-Pushnitski '06

$\{\nu_{k,q}\}_{k \in \mathbb{Z}_+}$ eig. of the s.a., compact Toeplitz operator $p_q \mathbb{1}_{\mathcal{O}} p_q$,
 p_q : proj. onto $\text{Ker}(H_{0,\perp} - \Lambda_q)$

$$\ln \nu_{k,q} = -k \ln k + (\mathfrak{C}(\mathcal{O}) - \ln 2) k + o(k), \quad k \rightarrow \infty,$$

$$\mathfrak{C}(\mathcal{O}) := \ln (eb \text{Cap}(\overline{\mathcal{O}})^2)$$

$\text{Cap}(K)$: *the logarithmic capacity* of K : $\text{Cap}(K) := e^{-\mathcal{I}(K)}$ where

$$\mathcal{I}(K) := \inf_{\mu \in \mathfrak{M}(K)} \int_{K \times K} \ln |x - y|^{-1} d\mu(x) d\mu(y).$$

Consequence:

Filonov-Pushnitski '06 ($H_{0,\perp} \pm \mathbb{1}_{\mathcal{O}}$),

Pu.-Roz '07/GoKaPe '14 ($H_{\perp,\pm}$: **Diri./Neu.** on $\partial\mathcal{O}$, $\mathcal{O} \subset \mathbb{R}^2$)



$\lambda_{k,q}$ eigenvalues of $H_{0,\perp} \pm \mathbb{1}_{\mathcal{O}}$ (or $H_{\perp,\pm}$) near Λ_q satisfy:

$$\ln(\pm(\lambda_{k,q} - \Lambda_q)) = -k \ln k + (\mathfrak{e}(\mathcal{O}) - \ln 2)k + o(k), \quad k \rightarrow \infty,$$

for $\mathcal{N}_{\pm}(\lambda) := \#\{\lambda_{k,q} \in \Lambda_q \pm (\lambda, r_0)\}$, $0 < \lambda < r_0 < 2b$, we have

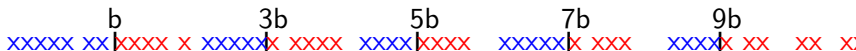
$$\mathcal{N}_{\pm}(\lambda) = \Phi_1(\lambda, \mathfrak{e}(\mathcal{O}) - \ln 2) + o\left(\frac{|\ln \lambda|}{\ln_2(\lambda)^2}\right) \quad \lambda \searrow 0$$

$$\Phi_1(\lambda; C) := \Phi_0(\lambda) \left(1 + \frac{\ln_3(\lambda)}{\ln_2(\lambda)} + \frac{C}{\ln_2(\lambda)}\right), \quad \ln_3(\lambda) := \ln \ln_2(\lambda),$$

Consequence:

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$\lambda_{k,q}$ eigenvalues of $H_{0,\perp} \pm \mathbf{1}_O$ (or $H_{\perp,\pm}$) near Λ_q satisfy:

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Asymptotic behavior of the SSF near Λ_q in the 3D case

Let \mathcal{O}_{in} be the projection of Ω_{in} onto the plane \perp to \vec{B} :

$$\mathcal{O}_{\text{in}} := \pi_{\perp}(\Omega_{\text{in}}) := \{X_{\perp} \in \mathbb{R}^2 \mid \exists x_{\parallel} \in \mathbb{R} \text{ s.t. } (X_{\perp}, x_{\parallel}) \in \Omega_{\text{in}}\}$$

Theorem 2

Let $q \in \mathbb{Z}_+$. Under assumption \mathcal{A} (see after) for Ω_{in} , we have

$$(i) \quad \xi(\Lambda_q - \lambda; H_-, H_0) = -\frac{1}{2} \Phi_1(\lambda; \mathfrak{C}(\mathcal{O}_{\text{in}})) + o\left(\frac{|\ln \lambda|}{\ln_2(\lambda)^2}\right)$$

$$(ii) \quad \xi(\Lambda_q + \lambda; H_{\pm}, H_0) = \pm \frac{1}{4} \Phi_1(\lambda; \mathfrak{C}(\mathcal{O}_{\text{in}})) + o\left(\frac{|\ln \lambda|}{\ln_2(\lambda)^2}\right)$$

as $\lambda \downarrow 0$.

$$\mathfrak{C}(\mathcal{O}) := \ln(\text{eb Cap}(\overline{\mathcal{O}})^2)$$

$$\Phi_1(\lambda; C) := \Phi_0(\lambda) \left(1 + \frac{\ln_3(\lambda)}{\ln_2(\lambda)} + \frac{C}{\ln_2(\lambda)}\right), \quad \Phi_0(\lambda) := \frac{|\ln \lambda|}{\ln_2(\lambda)}$$

$$\ln_3(\lambda) := \ln \ln_2(\lambda), \quad \ln_2(\lambda) := \ln \ln(\lambda),$$

Definition: Ω_{in} satisfies **assumption \mathcal{A}** if:

\exists **adjacent sequences** $\{\Omega_{j,<}\}_{j \in \mathbb{N}} \nearrow$ and $\{\Omega_{j,>}\}_{j \in \mathbb{N}} \searrow$ s.t.:

(i) $\bar{\Omega}_{j,<} \subset \Omega_{\text{in}} \subset \bar{\Omega}_{\text{in}} \subset \Omega_{j,>}$, $j \in \mathbb{N}$;

(ii) For $\mathcal{O}_{j,<} := \pi_{\perp}(\Omega_{j,<})$ and $\mathcal{O}_{j,>} := \pi_{\perp}(\Omega_{j,>})$, $j \in \mathbb{N}$,
 $\partial \mathcal{O}_{j,<}$ and $\partial \mathcal{O}_{j,>}$ are Lipschitz

(iii) For $\mathcal{O}_{\text{in}} := \pi_{\perp}(\Omega_{\text{in}})$;

$$\lim_{j \rightarrow \infty} \text{Cap}(\bar{\mathcal{O}}_{j,<}) = \lim_{j \rightarrow \infty} \text{Cap}(\bar{\mathcal{O}}_{j,>}) = \text{Cap}(\bar{\mathcal{O}}_{\text{in}}),$$

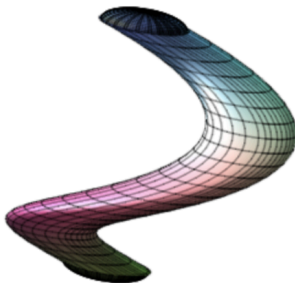
(iv) For any $j \in \mathbb{N}$,

$$\inf_{x_{\perp} \in \mathcal{O}_{j,<}} \int_{\mathbb{R}} \mathbf{1}_{\Omega_{j+1,<}}(X_{\perp}, x_{\parallel}) dx_{\parallel} > 0, \quad j \in \mathbb{N}.$$

Evidently, any ball (or ellipsoid) in \mathbb{R}^3 satisfies assumption \mathcal{A} .

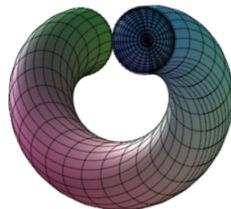
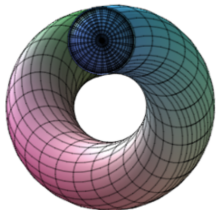
Exemples

Ω_{in}



\mathcal{O}_{in} smooth

\mathcal{O}_{in} not Lipschitz



III) Ideas of the proof

We have

$$\xi(E; H_{\pm}, H_0) = -\xi\left(\frac{1}{E}; H_{\pm}^{-1}, H_0^{-1}\right)$$

H_{\pm}^{-1} and H_0^{-1} defined in $L^2(\mathbb{R}^3)$ and $H_{\pm}^{-1} = H_0^{-1} \mp V_{\pm}$ with

$$V_+ := H_0^{-1} - H_+^{-1} \geq 0, \quad V_- := H_-^{-1} - H_0^{-1} \geq 0$$

Step 1: Pushnitski's representation formula (Push. '97)

\implies study of the Birman-Schwinger operator:

$$\begin{aligned} V_{\pm}^{\frac{1}{2}} \left(H_0^{-1} - E^{-1} \right)^{-1} V_{\pm}^{\frac{1}{2}} &= E V_{\pm}^{\frac{1}{2}} H_0 (E - H_0)^{-1} V_{\pm}^{\frac{1}{2}} \\ &= -E V_{\pm} - E^2 V_{\pm}^{\frac{1}{2}} (H_0 - E)^{-1} V_{\pm}^{\frac{1}{2}} \end{aligned}$$

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Step 2: Let $\mathcal{M}_q^\pm := \Lambda_q^2 V_\pm^{\frac{1}{2}}(p_q \otimes p_{\parallel})V_\pm^{\frac{1}{2}}$; $p_{\parallel} = \langle \cdot, 1 \rangle$.
Following Fe-Ra'04, as $\lambda \downarrow 0$, we have

Proposition 1

$$\begin{aligned} \xi(\Lambda_q - \lambda; H_+, H_0) &= O(1), \\ -\mathrm{Tr} \mathbf{1}_{[1, +\infty)} \left(\frac{\mathcal{M}_q^-}{(1-\epsilon)2\sqrt{\lambda}} \right) + O(1) &\leq \xi(\Lambda_q - \lambda; H_-, H_0) \\ &\leq -\mathrm{Tr} \mathbf{1}_{[1, +\infty)} \left(\frac{\mathcal{M}_q^-}{(1+\epsilon)2\sqrt{\lambda}} \right) + O(1) \\ \frac{1}{\pi} \mathrm{Tr} \arctan \left(\frac{\mathcal{M}_q^+}{(1+\epsilon)2\sqrt{\lambda}} \right) + O(1) &\leq \xi(\Lambda_q + \lambda; H_+, H_0) \\ &\leq \frac{1}{\pi} \mathrm{Tr} \arctan \left(\frac{\mathcal{M}_q^+}{(1-\epsilon)2\sqrt{\lambda}} \right) + O(1) \\ -\frac{1}{\pi} \mathrm{Tr} \arctan \left(\frac{\mathcal{M}_q^-}{(1-\epsilon)2\sqrt{\lambda}} \right) + O(1) &\leq \xi(\Lambda_q + \lambda; H_-, H_0) \\ &\leq -\frac{1}{\pi} \mathrm{Tr} \arctan \left(\frac{\mathcal{M}_q^-}{(1+\epsilon)2\sqrt{\lambda}} \right) + O(1) \end{aligned}$$

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The non zero eigenvalues of

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coincide with eigenvalues of

$$\Lambda_q^2 p_q \left(\int_{\mathbb{R}} V_\pm dx_3 \right) p_q$$

($0 \neq \text{e.v.}(TT^*) = \text{e.v.}(T^*T)$)

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Proposition 2

Let $\Omega_{<} \subset \Omega_{\text{in}} \subset \Omega_{>}$ as in Assumption \mathcal{A}

Let $\mathcal{E}_q(\Omega_{>}) = \{f \in L^2(\mathbb{R}^3) \cap C^\infty(\mathbb{R}^3); (H_0 - \Lambda_q)f = 0 \text{ on } \Omega_{>}\}$

Then there exists \mathcal{L}_q a finite codimension subspaces of $\mathcal{E}_q(\Omega_{>})$
and $C > 1$ s. t. $\forall f \in \mathcal{L}_q$,

$$\frac{1}{C} \langle f, \mathbf{1}_{\Omega_{<}} f \rangle_{L^2(\mathbb{R}^3)} \leq \langle H_0 w f, V_{\pm} H_0 w f \rangle_{L^2(\mathbb{R}^3)} \leq C \langle f, \mathbf{1}_{\Omega_{>}} f \rangle_{L^2(\mathbb{R}^3)}$$

\implies Theorem follows using that

$$c \mathbf{1}_{\mathcal{O}_{j,<}}(X_{\perp}) \leq \int_{\mathbb{R}} \mathbf{1}_{\Omega_{j+1,<}}(X_{\perp}, x_{\parallel}) dx_{\parallel}$$
$$\int_{\mathbb{R}} \mathbf{1}_{\Omega_{j+1,>}}(X_{\perp}, x_{\parallel}) dx_{\parallel} \leq C \mathbf{1}_{\mathcal{O}_{j+1,>}}(X_{\perp})$$

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Proof of the lower bound of Proposition 2

$$\langle H_0 w f, V_{\pm} H_0 w f \rangle_{L^2(\mathbb{R}^3)} \geq \frac{1}{C} \langle f, \mathbf{1}_{\Omega_{<}} f \rangle_{L^2(\mathbb{R}^3)}$$

For V_+ :

$$V_+ := H_0^{-1} - H_{+,in}^{-1} \oplus H_{+,ex}^{-1} = V_{+,0} - H_{+,in}^{-1} \oplus 0, \text{ with}$$

$$V_{+,0} := H_0^{-1} - 0 \oplus H_{+,ex}^{-1} \geq \left(H_0^{-1} - (H_0 + \mathbf{1}_{\Omega_{<}})^{-1} \right)$$

because

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We conclude using

$$H_0^{-1} - (H_0 + \mathbf{1}_{\Omega_{<}})^{-1} =$$

$$H_0^{-1} \mathbf{1}_{\Omega_{<}} \left(I - \mathbf{1}_{\Omega_{<}} (H_0 + \mathbf{1}_{\Omega_{<}})^{-1} \mathbf{1}_{\Omega_{<}} \right) \mathbf{1}_{\Omega_{<}} H_0^{-1} \geq H_0^{-1} \mathbf{1}_{\Omega_{<}} H_0^{-1}$$

on a finite codim. subsp.

Proof of the lower bound of Proposition 2

$$\langle H_0 w f, V_{\pm} H_0 w f \rangle_{L^2(\mathbb{R}^3)} \geq \frac{1}{C} \langle f, \mathbf{1}_{\Omega_{<}} f \rangle_{L^2(\mathbb{R}^3)}$$

For V_+ :

$$V_+ := H_0^{-1} - H_{+,in}^{-1} \oplus H_{+,ex}^{-1} = V_{+,0} - H_{+,in}^{-1} \oplus 0, \text{ with}$$

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$$V_- := H_{-,in}^{-1} \oplus H_{-,ex}^{-1} - H_0^{-1} = V_{-,0} - \delta H_{-,in}^{-1} \oplus 0, \quad \delta > 0 \text{ with}$$

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There exists $\kappa > 0$ s.t.:

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Conclusion on the 3D magnetic problem with obstacle Ω_{in} :

As $\lambda \downarrow 0$, the SSF satisfies in particular :

$$\xi(\Lambda_q + \lambda; H_{\pm, \text{ex}}, H_0) = \pm \frac{1}{4} \left(\frac{|\ln \lambda|}{\ln_2(\lambda)} + \frac{|\ln \lambda| \ln_3(\lambda)}{\ln_2(\lambda)^2} + \frac{|\ln \lambda| \mathfrak{C}(\mathcal{O}_{\text{in}})}{\ln_2(\lambda)^2} \right) + o\left(\frac{|\ln \lambda|}{\ln_2(\lambda)^2}\right)$$

with:

$$\mathfrak{C}(\mathcal{O}) := \ln(\text{eb Cap}(\overline{\mathcal{O}})^2)$$

$$\mathcal{O}_{\text{in}} := \pi_{\perp}(\Omega_{\text{in}}) := \{X_{\perp} \in \mathbb{R}^2 \mid \exists x_{\parallel} \in \mathbb{R} \text{ s.t. } (X_{\perp}, x_{\parallel}) \in \Omega_{\text{in}}\}$$

Under assumption \mathcal{A}

Question: Is it a technical assumption? Or more fundamental?

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